



University  
of Glasgow

# Instabilities in a Planar-Aligned Active Liquid Crystal

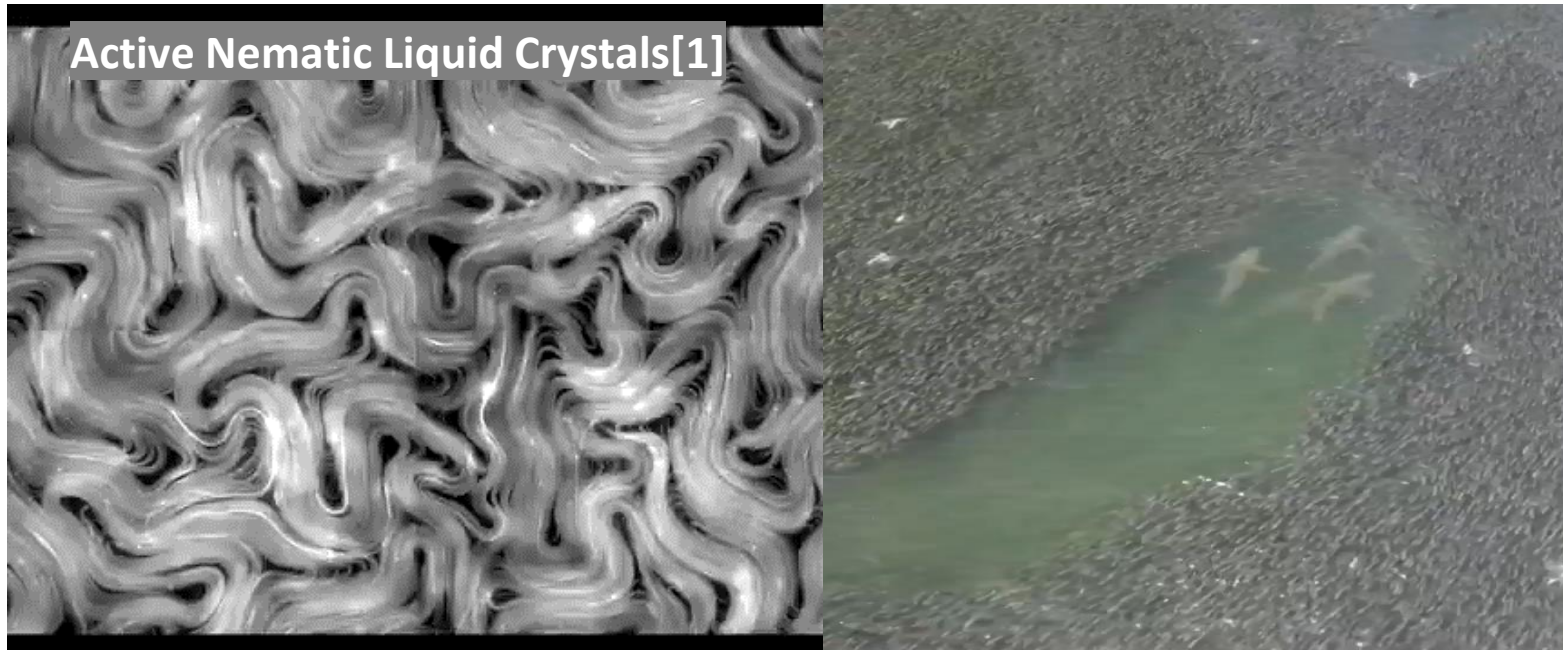
Ijuptil Joseph Kwajighu  
(2592566j@student.gla.ac.uk)

School of Mathematics and Statistics, University of Glasgow

Supervisors: Prof. Nigel J. Mottram, Dr Katarzyna N. Kowal, Dr Joseph R. L. Cousins

# Introduction

- ❖ Liquid crystals are intermediate phases, between a solid crystal and an isotropic liquid.
- ❖ Active liquid crystals are made of objects that can generate flow.
- ❖ These systems swim in patterns that suggest **long-range collective ordering**.
- ❖ The **activity** in these systems relies on **continuous energy production** by the individual particles.

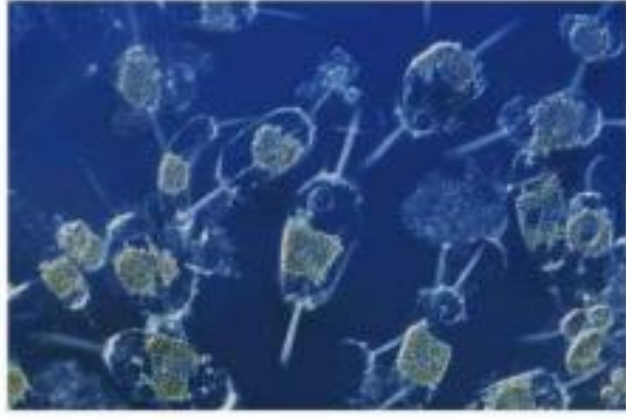
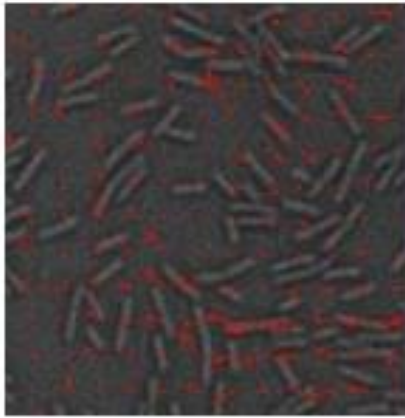


1. S. J. DeCamp, G. S. Redner, A. Baskaran, M. F. Hagan, and Z. Dogic, *Nat Mater* 14, 1110, 2015.



# Introduction

## Examples of Active Liquid Crystals



bacteria

phytoplankton

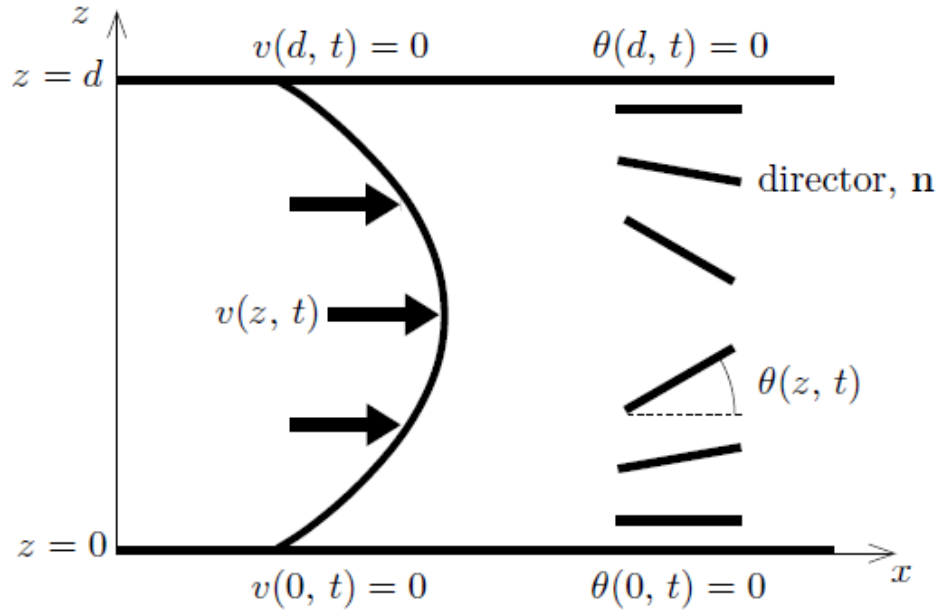
zooplankton

fish

- ❖ Will an orienting field increase or decrease the flow generated by active nematics?
- ❖ The use of active liquid crystals in designing sensors is an interesting possible application.
- ❖ We have used mathematical modelling of active nematics using an adapted version of the continuum theory of nematics.

# Activity-Driven Channel Flow

## Model geometry



An active nematic between two solid plates at  $z=0, d$  with flow in the  $x$ -direction [2].

- ❖ The configuration and organization of the active agents are deformed by flow, and the active agents also generate flow.

## Ericksen-Leslie Equations for an Active Nematic

- ❖ The simplest model uses the Ericksen-Leslie theory with just one extra term in the stress tensor [3]

$$\boxed{\gamma_1 \theta_t} = \underbrace{(K_1 \cos^2 \theta + K_3 \sin^2 \theta) \theta_{zz} + (K_3 - K_1) \sin \theta \cos \theta (\theta_z)^2}_{\text{elastic terms}} - \underbrace{m(\theta) v_z}_{\text{director-flow coupling}},$$

$$\boxed{\rho v_t} = \underbrace{(g(\theta) v_z)}_{\text{flow inertia}} + \underbrace{m(\theta) \theta_t}_{\text{fluid viscosity}} + \underbrace{\zeta \cos \theta \sin \theta}_z \underbrace{\zeta \cos \theta \sin \theta}_{\text{activity}}$$

- ❖ The stress tensor is written as:  $\tau = \tau_{nem} + \underbrace{\zeta (\mathbf{n} \otimes \mathbf{n})}_{\text{activity}}$

1. F. M. Leslie. Continuum theory for nematic liquid crystals. *Continuum Mechanics and Thermodynamics*, 4(3):167–175, 1992.  
 2. J. Walton, G. McKay, M. Grinfeld, and N. J. Mottram. *Mathematical modelling of active nematic liquid crystals in coned regions*. PhD thesis, University of Strathclyde, 2020



## ❖ Linearisation around a planar alignment

$$\gamma_1 \frac{\partial \theta}{\partial t} - K_1 \frac{\partial^2 \theta}{\partial z^2} - \alpha_3 \frac{\partial v}{\partial z} = 0, \quad \eta_1 \frac{\partial^2 v}{\partial z^2} + \alpha_3 \frac{\partial^2 \theta}{\partial z \partial t} + \zeta \frac{\partial \theta}{\partial z} = 0.$$

## ❖ The director angle and flow velocity

$$\theta(z, t) = \theta_0 \left[ \cos \left( q - \frac{2q}{d} z \right) - \cos(q) \right] \exp \left( -\frac{t}{\tau} \right),$$

$$v(z, t) = v_0 \left[ d \sin \left( q - \frac{2q}{d} z \right) + (2z - d) \sin(q) \right] \exp \left( -\frac{t}{\tau} \right).$$

**Time constant**

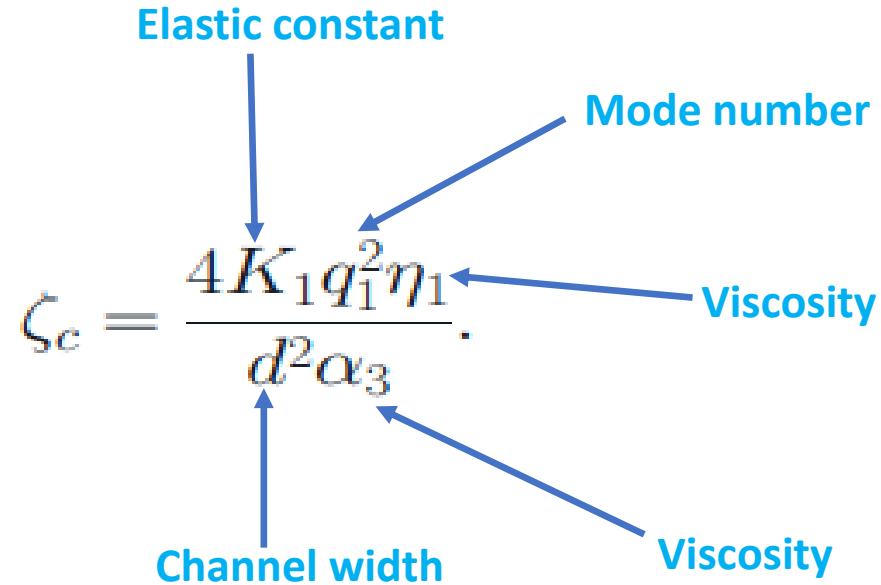
**Mode number**

$$\tau = \alpha \gamma_1 \left( \frac{4K_1 q^2}{d^2} - \frac{\alpha_3 \zeta}{\eta_1} \right)^{-1} \quad r(q) = q - (1 - \alpha) \tan q + \frac{\alpha_2 \xi \alpha}{\eta_2} \left( \frac{4K_3 q^2}{d^2} + \frac{\alpha_2 \xi}{\eta_2} \right)^{-1} \tan q = 0.$$

❖ When the viscosity  $\alpha_3 < 0$  we can see that  $\tau > 0$  whenever  $\zeta > 0$ , so that the system is **stable**.

❖ When the viscosity  $\alpha_3 > 0$  we can see that  $\tau > 0$  whenever  $\zeta < 0$ , so that the system is **stable**.

❖ **Instability:** The critical activity occurs at the **first mode**,  $n = 1$ , which leads to instability when  $\zeta = \zeta_c$

$$\zeta_c = \frac{4K_1 q_1^2 \eta_1}{d^2 \alpha_3}$$


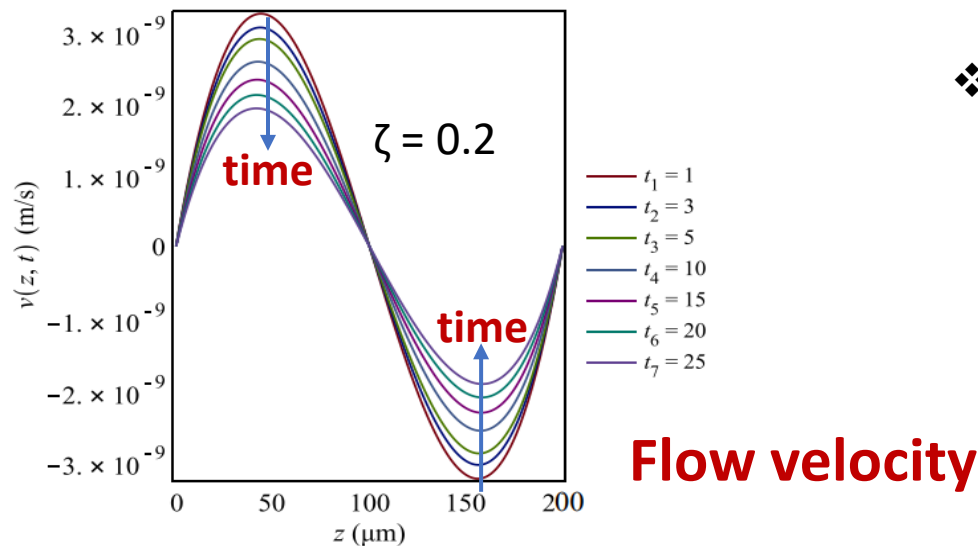
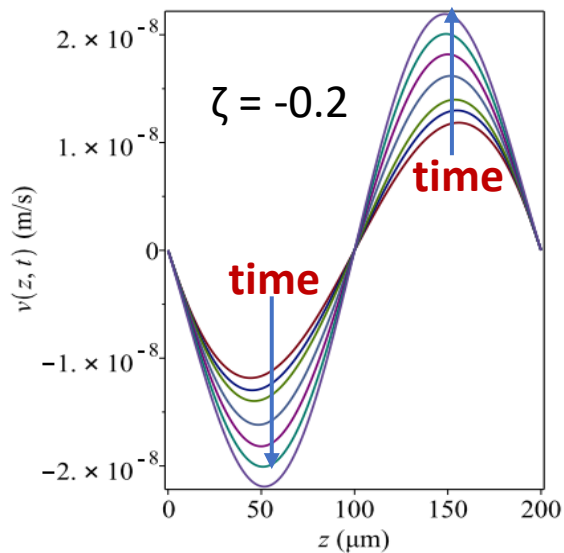
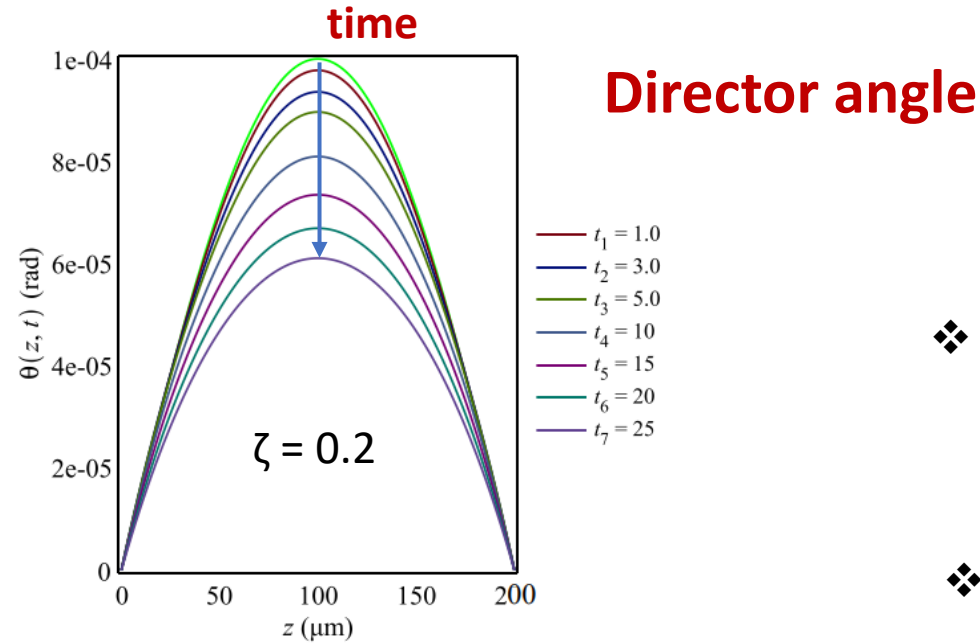
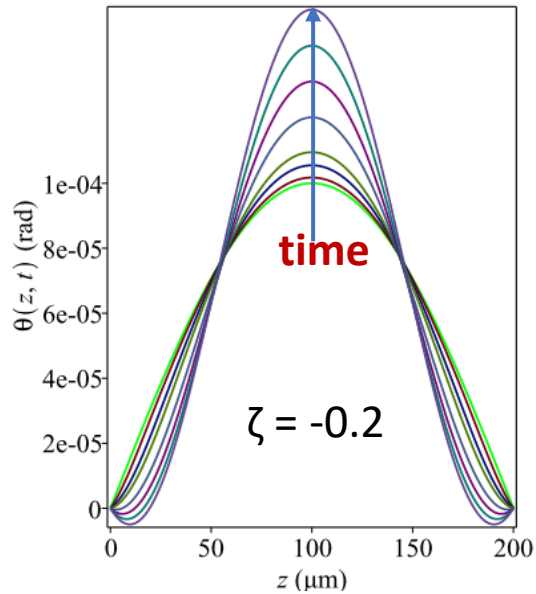
The diagram illustrates the components of the critical activity equation  $\zeta_c = \frac{4K_1 q_1^2 \eta_1}{d^2 \alpha_3}$ . Blue arrows point from labels to the corresponding terms in the equation:

- Elastic constant** points to  $K_1$ .
- Mode number** points to  $q_1$ .
- Viscosity** points to  $\eta_1$ .
- Channel width** points to  $d$ .
- Viscosity** points to  $\alpha_3$ .

❖ When the active agent is rod-like then  $\alpha_3 < 0$  and  $\zeta_c < 0$ .

❖ When the active agent is disc-like then  $\alpha_3 > 0$  and  $\zeta_c > 0$ .

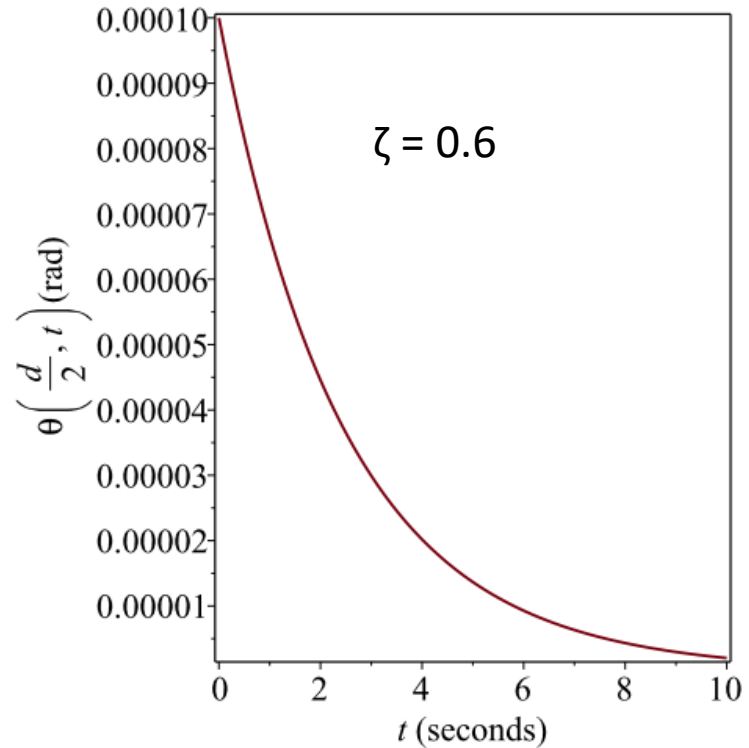
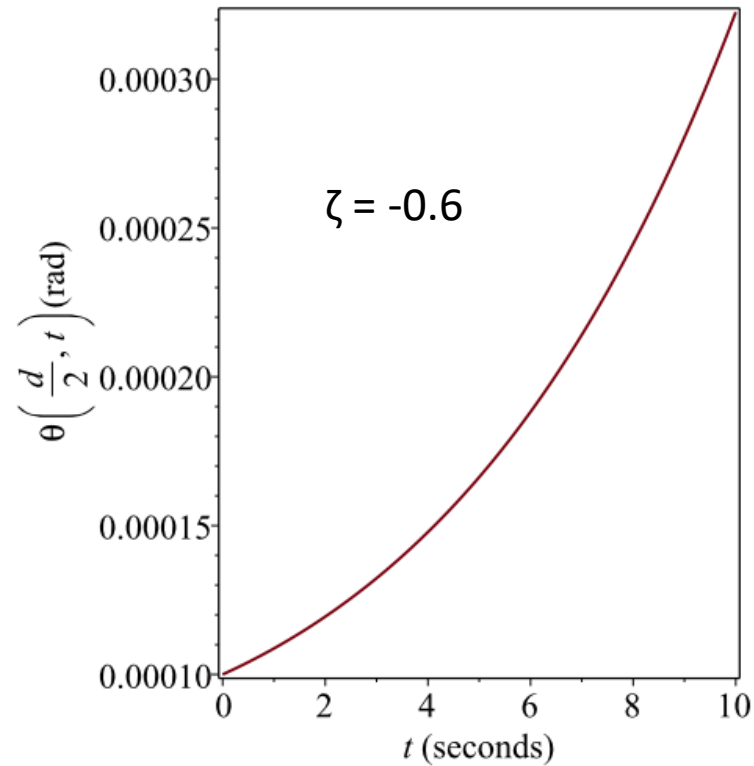
# Results



- ❖ For the **negative activity** parameter the **flow grows** with time, for the **positive activity** parameter, the **flow decays**.
- ❖ As the **magnitude of the activity** increases, the active system tends to **generate flow**.
- ❖ The **fastest-growing mode** for each value of the activity is the mode with the **most negative real part** of the time constant.

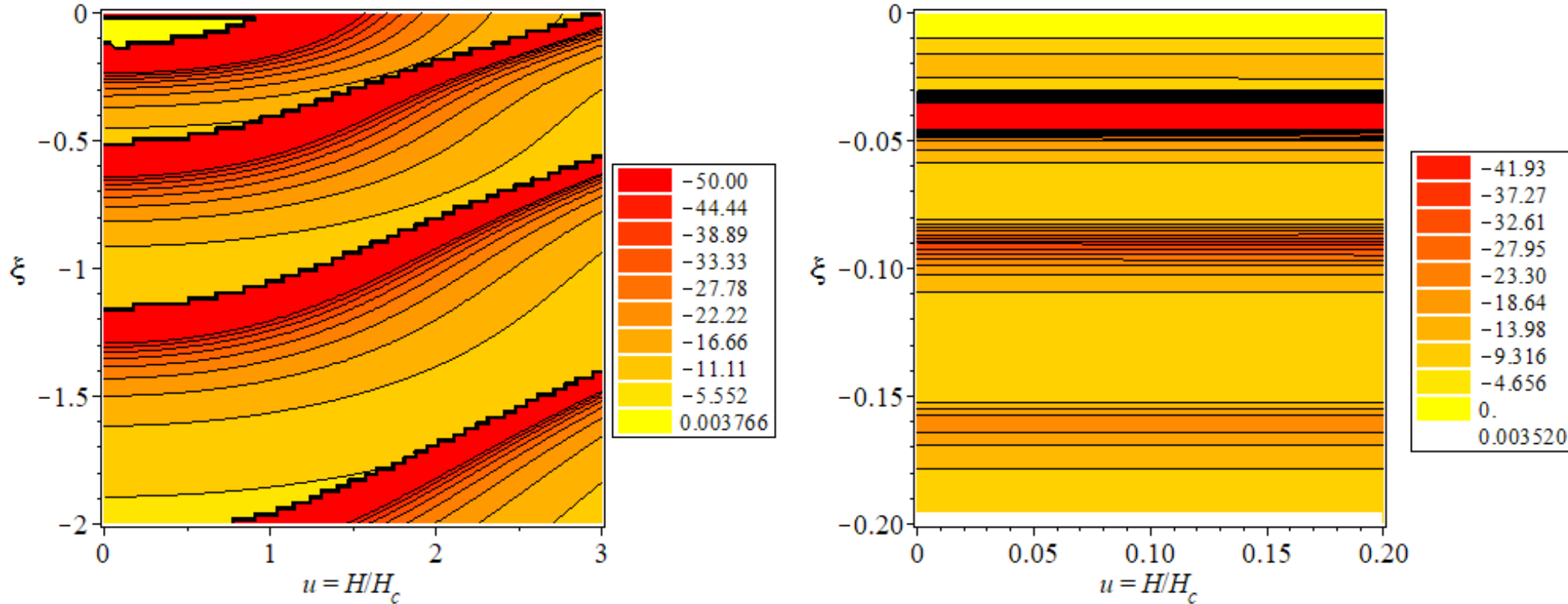


The director angle in the middle of the layer as a function of  $t$



❖ The activity coefficient controls the generation and growth/decay of flow in the system.

- ❖ Contour plots of the minimum values of the **time constant**  $\tau$  and the corresponding imaginary values.



- ❖ In future research, we will solve the full coupled nonlinear PDEs using an appropriate numerical methods.